# Derivatives of Bivalent Germanium, Tin, and Lead. Part 21.1 Tin(II) Formate: A Reinvestigation 

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#### Abstract

The structure of the title compound has been determined from diffractometer data using Patterson and Fourier methods. Crystals are triclinic, $P \overline{1}$, with unit-cell dimensions $a=5.303(2), b=8.834(3), c=5.271(2) \AA, \alpha=$ 104.38(2), $\beta=105.75(2)$, and $\gamma=82.82(2)^{\circ}$. Refinement has been carried by full-matrix least squares to a final $R$ value of 0.0503 for 1309 independent non-zero reflections. Crystals consist of infinite two-dimensional sheets in which all the formate groups bridge adjacent tin atoms. Tin forms two short [2.14(1) and 2.20(1) A] and two long [2.36(1) $\AA$ ] bonds in a pseudo-trigonal-bipyramidal arrangement with the tin lone pair occupying the remaining equatorial site. The structure is disordered such that the formate groups occupy equally two sets of positions between the tin atoms. Vibrational and tin- 119 m Mössbauer spectra have been recorded, and together with the unit-cell data are compared with data reported previously.


Although tin(iI) carboxylates have been known for well over a century, ${ }^{2}$ and their synthesis and spectroscopic properties well documented, ${ }^{3}$ the structures of only three sophisticated and atypical examples have been determined: (i) dihydrogen ethylenediaminetetra-acetatostannate(II), in which the tin atom is centred in the rectangular face of a distorted triangular prism and coordinated by two nitrogen and four oxygen atoms; ${ }^{4}$ (ii) ditin(II) ethylenediaminetetra-acetate dihydrate, in which one tin atom experiences pentagonal-bipyramidal pseudo-seven-co-ordination with the lone pair occupying an equatorial site, whilst the second tin experiences irregular pseudo-eight-co-ordinate geometry in which the lone pair is again stereochemically active; ${ }^{5}$ and (iii) $\left[\mathrm{Sn}^{\mathrm{II}} \mathrm{Sn}^{\mathrm{IV}}\left(\mathrm{O}_{2} \mathrm{CC}_{6} \mathrm{H}_{4} \mathrm{NO}_{2}-o\right)_{4} \mathrm{O} \cdot \mathrm{C}_{4} \mathrm{H}_{8} \mathrm{O}\right]_{2}$, the tin(II) atoms of which enjoy pentagonal-pyramidal co-ordination, with the lone pair occupying the second apical site of a bipyramid. ${ }^{6}$ However, the structures have been determined of the two simple carboxylatostannate(II) anions $\left[\mathrm{Sn}\left(\mathrm{O}_{2} \mathrm{CH}\right)_{3}\right]^{-}\left(\right.$ref. 7 ) and $\left[\mathrm{Sn}\left(\mathrm{C}_{2} \mathrm{O}_{4}\right)_{2}\right]^{2-}$ (ref. 8) as their alkali-metal salts. The former anion has pyramidal geometry with unidentate formate groups, ${ }^{7}$ whilst in the latter both oxalate groups chelate the tin resulting in a distorted pseudo-trigonal-bipyramidal stereochemistry, with the lone pair in one of the equatorial sites. ${ }^{8}$

Donaldson and his co-workers ${ }^{3,9-11}$ have investigated in some detail the preparation and properties of the simplest tin(II) carboxylate, tin(II) formate, and have concluded from both tin-119m Mössbauer ${ }^{3}$ and i.r. data ${ }^{11}$ that this derivative enjoys pyramidal co-ordination for the tin atom. In this paper we report the results of an $X$-ray diffraction study of this compound, which shows that the proposed structure based on spectroscopic data is erroneous, together with a reinvestigation of its tin119 m Mossbauer, and vibrational, data.

[^0]
## EXPERIMENTAL

The crystalline sample of anhydrous tin(II) formate was prepared from $\operatorname{tin}(\mathrm{II})$ oxide and formic acid according to Donaldson and Knifton ${ }^{9}$ [Found: C, 11.3; H, 0.94. $\mathrm{Sn}\left(\mathrm{O}_{2} \mathrm{CH}\right)_{2}$ requires C, $11.5 ; \mathrm{H}, 0.97 \%$ ]. The white crystals (yield $81 \%$ ) started to darken at $156^{\circ} \mathrm{C}$, were black at $200^{\circ} \mathrm{C}$, but did not change shape on heating to $300^{\circ} \mathrm{C}$. A specimen of dimensions $c a .0 .4 \times 0.4 \times 0.3 \mathrm{~mm}$ was mounted under nitrogen in a Lindemann capillary and used for both the initial photography and also the subsequent intensity-data collection.

Table 1
Crystal data for $\operatorname{tin}$ (II) formate, $\mathrm{C}_{2} \mathrm{H}_{2} \mathrm{O}_{4} \mathrm{Sn}, M=208.72$

| Space group | This work Triclinic, $P \mathbf{I}$ | Ref. 9 Triclinic ${ }^{a}$ |
| :---: | :---: | :---: |
| $a / \AA$ | $5.303(2)$ | 10.56(20) |
|  |  | ( $\equiv 2 \times 5.28)$ |
| b/A | 8.834(3) | $8.81(20)$ |
| $c / \AA$ | 5.271 (2) | 10.56(20) |
|  |  | ( $\equiv 2 \times 5.28$ ) |
| $\alpha /^{\circ}$ | 104.38(2) | 82.83(8) |
|  | $(\equiv 90+14.38)$ | $\left(\equiv 90-\frac{1}{2} \times 14.38\right)$ |
| $\beta 1^{\circ}$ | 105.75(2) | $\begin{gathered} 74.83(8) \\ (\equiv 180-105.17) \end{gathered}$ |
| $\gamma{ }^{\circ}$ | 82.82(2) | 82.83(8) |
| $U / \AA^{3}$ | 229.8 | 937.6 |
| $D_{\text {c }} / \mathrm{g} \mathrm{cm}^{-3}$ | 3.02 | 2.96 |
| $D_{\mathrm{m}} / \mathrm{g} \mathrm{cm}^{-3}$ | $3.06{ }^{\text {b }}$ | $2.98{ }^{\circ}$ |
| $Z$ | 2 | 8 |
| $F(000)$ | 192 | 768 |
| Radiation | $\mathrm{Mo}-K_{\alpha}$ | $\mathrm{Cu}-K_{\alpha}$ |
| $\lambda / \AA$ | 0.7107 |  |
| $\mu\left(\mathrm{Mo}-K_{\alpha}\right) / \mathrm{cm}^{-1}$ | 54.62 |  |

${ }^{a}$ The $b$ and $c$ axes of the original cell have been interchanged.
${ }^{b}$ By flotation. ${ }^{c}$ By displacement.
Crystal Data.-Crystal data for $\operatorname{tin}(\mathrm{II})$ formate obtained in the present study are listed in Table 1 together with those obtained previously. ${ }^{9}$

The space group and initial cell parameters were deter-

[^1]mined from oscillation and zero- and first-layer Weissenberg photographs obtained on an equi-inclination Weissenberg camera. Relative intensities of 1309 reflections were collected up to $\theta 30.0^{\circ}$ with $\mathrm{Mo}-K_{\alpha}$ radiation by use of a Hilger and Watts Y290 single-crystal four-circle diffractometer, accurate cell dimensions being obtained by leastsquares refinement of ca. 20 reflections. Each reflection was counted for 30 s and the two associated background counts for 10 s . The reflections were brought to the same relative intensities by use of reference reflections taken for every one hundred reflections recorded; the diffractometer orientation was referenced every two hundred reflections. The intensities were corrected for Lorentz and polarisation effects, but not for absorption.

Structure Determination and Refinement.-The positional parameters of the tin atom were obtained from a threedimensional Patterson synthesis, and these were used to phase the initial structure-factor calculation. Since a subsequent Fourier synthesis did not yield any chemically sensible information, the position of the tin atom was refined by isotropic least squares, and a second Fourier synthesis afforded the positions of eight light atoms around the tin. These were assigned as half-weight oxygen atoms, and the refinement continued with the tin atom being allowed to vary anisotropically. After three cycles of refinement, a third Fourier synthesis yielded the positions of the two halfweight carbon atoms. After two more cycles of full-matrix refinement with only the tin atom varying anisotropically, followed by two cycles of full-matrix refinement with all the atoms varying anisotropically, the weighting scheme $w^{\frac{1}{2}}=1$ for $F_{0} \leqslant 27.0$ and $w^{\frac{1}{2}}=27 / F_{0}$ for $F_{0}>27.0$ was applied. Two additional cycles of full-matrix anisotropic refinement produced convergence at an ' $R$ ' value of 0.0503 . The scattering factors used were those for neutral atoms. ${ }^{12}$ Calculations were performed using the CRYSTALS suite of programs. ${ }^{13}$ The results are collected in Tables $1-3$.

Table 2
Final fractional atomic co-ordinates, with estimated standard deviations in parentheses

| Atom | Occupancy | $x / a$ | $y / b$ | $z / c$ |
| :---: | :---: | :---: | :---: | :---: |
| $\mathrm{Sn}(1)$ | 1.0 | $0.23635(7)$ | $0.13036(4)$ | $0.30167(7)$ |
| $\mathrm{O}(1)$ | 0.5 | $0.148(2)$ | $0.183(1)$ | 0.732(2) |
| $\mathrm{O}(2)$ | 0.5 | $0.308(2)$ | $0.180(1)$ | 0.908(2) |
| $\mathrm{O}(3)$ | 0.5 | $0.956(2)$ | 0.319(1) | 0.228(2) |
| $\mathrm{O}(4)$ | 0.5 | 0.477(2) | 0.323(1) | 0.517(2) |
| $\mathrm{O}(5)$ | 0.5 | $0.821(2)$ | 0.185(1) | $0.401(2)$ |
| $\mathrm{O}(6)$ | 0.5 | $0.644(2)$ | $0.183(1)$ | $0.248(2)$ |
| $\mathrm{O}(7)$ | 0.5 | 0.080 (2) | $0.330(1)$ | 0.113(2) |
| $\mathrm{O}(8)$ | 0.5 | 0.355 (2) | $0.329(1)$ | $0.645(2)$ |
| $\mathrm{C}(1)$ | 0.5 | 0.074(3) | 0.308(2) | 0.862(3) |
| $\mathrm{C}(2)$ | 0.5 | 0.358(3) | 0.308(2) | 0.874(3) |
| $\mathrm{C}(3)$ | 0.5 | $0.722(3)$ | $0.307(2)$ | 0.224(3) |
| $\mathrm{C}(4)$ | 0.5 | 0.717(3) | 0.307(2) | $0.511(3)$ |

Atomic numbering is shown in Figure 2. Final observed and calculated structure factors, and anisotropic thermal parameters, are listed in Supplementary Publication No. SUP 22308 ( 15 pp. ).*

Infrared spectra in the range $200-4000 \mathrm{~cm}^{-1}$ were recorded as Nujol and halogenocarbon mulls using a Perkin-Elmer 577 spectrometer. Reflectance spectra were

* For details see Notices to Authors No. 7, J.C.S. Dalton, 1977, Index issue.
12 'International Tables for $X$-Ray Crystallography,' Kynoch Press, Birmingham, 1962, vol. 3.
${ }^{13}$ C. M. Carruthers, personal communication.
obtained from a pressed disc in the range $40-400 \mathrm{~cm}^{-1}$ using a Beckmann R11C F.S. 720 Fourier-transform spectrophotometer, Raman spectra using a Carey 81 spectrometer

Table 3
Bond distances $(\AA)$ and angles $\left(^{\circ}\right)$ in $\operatorname{tin}(I I)$ formate (with estimated standard deviations in parentheses)
(a) Bond distances

| $\mathrm{Sn}(1)-\mathrm{O}(1)$ | $2.36(1)$ | $\mathrm{Sn}(1)-\mathrm{O}(5)$ | $2.36(1)$ |
| :--- | :--- | :--- | :--- |
| $\mathrm{Sn}(1)-\mathrm{O}(2)$ | $2.36(1)$ | $\mathrm{Sn}(1)-\mathrm{O}(6)$ | $2.37(1)$ |
| $\mathrm{Sn}(1)-\mathrm{O}(3)$ | $2.13(1)$ | $\mathrm{Sn}(1)-\mathrm{O}(7)$ | $2.20(1)$ |
| $\mathrm{Sn}(1)-\mathrm{O}(4)$ | $2.14_{5}(1)$ | $\mathrm{Sn}(1)-\mathrm{O}(8)$ | $2.20(1)$ |
| $\mathrm{C}(1)-\mathrm{O}(1)$ | $1.23(2)$ | $\mathrm{C}(3)-\mathrm{O}(3)$ | $1.25(2)$ |
| $\mathrm{C}(1)-\mathrm{O}(7)$ | $1.28(2)$ | $\mathrm{C}(3)-\mathrm{O}(6)$ | $1.26(2)$ |
| $\mathrm{C}(2)-\mathrm{O}(2)$ | $1.26(2)$ | $\mathrm{C}(4)-\mathrm{O}(4)$ | $1.275(2)$ |
| $\mathrm{C}(2)-\mathrm{O}(8)$ | $1.26(2)$ | $\mathrm{C}(4)-\mathrm{O}(5)$ | $1.23(2)$ |
| $(b) \operatorname{Bond}$ angles |  |  |  |
| $\mathrm{O}(1)-\mathrm{Sn}(1)-\mathrm{O}(2)$ | $158.6(4)$ | $\mathrm{O}(5)-\mathrm{Sn}(1)-\mathrm{O}(6)$ | $157.7(4)$ |
| $\mathrm{O}(3)-\mathrm{Sn}(1)-\mathrm{O}(4)$ | $80.7(4)$ | $\mathrm{O}(7)-\mathrm{Sn}(1)-\mathrm{O}(8)$ | $78.7(4)$ |
| $\mathrm{O}(1)-\mathrm{Sn}(1)-\mathrm{O}(3)$ | $85.7(4)$ | $\mathrm{O}(5)-\mathrm{Sn}(1)-\mathrm{O}(7)$ | $87.6(4)$ |
| $\mathrm{O}(1)-\mathrm{Sn}(1)-\mathrm{O}(4)$ | $78.45(4)$ | $O(5)-\mathrm{Sn}(1)-\mathrm{O}(8)$ | $84.4(4)$ |
| $\mathrm{O}(2)-\mathrm{Sn}(1)-\mathrm{O}(3)$ | $77.6(4)$ | $\mathrm{O}(6)-\mathrm{Sn}(1)-\mathrm{O}(7)$ | $84.5(4)$ |
| $\mathrm{O}(2)-\mathrm{Sn}(1)-\mathrm{O}(4)$ | $85.7(4)$ | $\mathrm{O}(6)-\mathrm{Sn}(1)-\mathrm{O}(8)$ | $78.05(4)$ |
| $\mathrm{Sn}(1)-\mathrm{O}(1)-\mathrm{C}(1)$ | $128.2(9)$ | $\mathrm{Sn}(1)-\mathrm{O}(5)-\mathrm{C}(4)$ | $131.3(9)$ |
| $\mathrm{Sn}(1)-\mathrm{O}(2)-\mathrm{C}(2)$ | $127.5(9)$ | $\mathrm{Sn}(1)-\mathrm{O}(6)-\mathrm{C}(3)$ | $131.0(9)$ |
| $\mathrm{Sn}(1)-\mathrm{O}(3)-\mathrm{C}(3)$ | $118.7(9)$ | $\mathrm{Sn}(1)-\mathrm{O}(7)-\mathrm{C}(1)$ | $116.2(9)$ |
| $\mathrm{Sn}(1)-\mathrm{O}(4)-\mathrm{C}(4)$ | $117.6_{5}(9)$ | $\mathrm{Sn}(1)-\mathrm{O}(8)-\mathrm{C}(2)$ | $115.7(9)$ |
| $\mathrm{O}(1)-\mathrm{C}(1)-\mathrm{O}(7)$ | $122.35(13)$ | $\mathrm{O}(3)-\mathrm{C}(3)-\mathrm{O}(6)$ | $122.5(12)$ |
| $\mathrm{O}(2)-\mathrm{C}(2)-\mathrm{O}(8)$ | $122.7(12)$ | $\mathrm{O}(4)-\mathrm{C}(4)-\mathrm{O}(5)$ | $123.7(12)$ |

equipped with an argon-ion laser (514.9-nm exciting line). Tin- 119 m Mössbauer data were collected at 77 K and ambient temperature (using equipment described previously) against a $\mathrm{Ba}^{119 \mathrm{~m}} \mathrm{SnO}_{3}$ source (Radiochemical Centre, Amersham). Data reduction and curve-fitting to Lorentzian line shapes was accomplished by usual least-squares methods.

## DISCUSSION

Tin-119m Mössbauer and vibrational spectroscopic data previously reported for simple tin(II) carboxylates have been interpreted in terms of trigonal-pyramidal coordination for the tin(II) atoms with bridging and/or chelating carboxylate groups. ${ }^{3,10,11}$ As the results of the present $X$-ray diffraction study show, this inference is not valid, at least in the case of $\operatorname{tin}$ (II) formate.

Crystals of $\operatorname{tin}$ (II) formate are composed of infinite twodimensional sheets in which tin(II) atoms are bridged by formate groups (Figure 1). Each tin(II) atom is bonded to four oxygen atoms from four different formate groups, the resulting 'local' stereochemistry at tin being that of a distorted pseudo-trigonal bipyramid, with the lone pair of electrons of the tin(iI) atom located in one of the equatorial sites. The geometry is thus similar to that previously observed in bis(1-phenylbutane-1,3-dionato)$\operatorname{tin}(\mathrm{II}),{ }^{14} \mathrm{Sn}\left(\mathrm{SO}_{4}\right) \cdot 2 \mathrm{SC}\left(\mathrm{NH}_{2}\right)_{2},{ }^{15} \mathrm{SnCl}_{2} \cdot$ dioxan, ${ }^{16}$ and the $\left[\mathrm{Sn}\left(\mathrm{O}_{4} \mathrm{C}_{2}\right)_{2}\right]^{2-}$ anion. ${ }^{8}$ The structure of $\operatorname{tin}($ II $)$ formate is, however, disordered such that the formate groups occupy equally two sets of positions between adjacent tin atoms, thus forming two superimposed independent twodimensional networks (Figure 2). Within each network, the repeating unit involves four mutually bridged tin
${ }^{14}$ P. F. R. Ewings, P. G. Harrison, and T. J. King, J.C.S. Dalton, 1975, 1455.

15 J. D. Donaldson, D. G. Nicholson, D. C. Puxley, and R. A. Howie, J.C.S. Dalton, 1973, 1810.
${ }^{16}$ E. Hough and D. G. Nicholson, J.C.S. Dalton, 1976, 1631.
atoms, which is consistent with a unit cell having $a$ and $c$ dimensions exactly twice those observed here as obtained by Donaldson and Knifton ${ }^{9}$ in their study, suggesting


Figure 1 Projection on to the ac plane showing the formate bridging within the two-dimensional sheets. Only one of the two disordered networks is shown. The broken lines refer to symmetry-related atoms in the lower layer
that these workers were investigating a crystal which was not disordered. The effect of the disorder observed in the present study is such as to make the overall coordination about each tin atom identical, since about each tin atom are half-weight contributions from two sets of four formate groups. The repeating unit along both the $a$ and $c$ directions therefore now involves only one tin


Figure 2 View of the structure showing the two disordered networks (solid and broken lines, respectively) along the bisector of the $\mathrm{O}_{\mathrm{ax}}-\mathrm{Sn}-\mathrm{O}_{\mathrm{ax}}$. and $\mathrm{O}_{\text {eq }}-\mathrm{Sn}-\mathrm{O}_{\text {eq. }}$. angles, as well as the atomic numbering
atom, although equivalent oxygen positions about adjacent tin atoms belong to different networks. Thus, referring to Figure 2, whereas $\mathrm{O}(5), \mathrm{O}(6), \mathrm{O}(7)$, and $\mathrm{O}(8)$ about tin atom $\mathrm{Sn}^{\mathrm{I}}$ belong to network A (full lines) and $\mathrm{O}(1), \mathrm{O}(2), \mathrm{O}(3)$, and $\mathrm{O}(4)$ about $\mathrm{Sn}^{\mathrm{I}}$ belong to network B (broken lines), the situation is reversed for the adjacent tin atoms $\mathrm{Sn}^{\mathrm{II}}, \mathrm{Sn}^{\mathrm{III}}, \mathrm{Sn}^{\mathrm{IV}}$, and $\mathrm{Sn}^{\mathrm{V}}$. The two-dimensional sheets are stacked parallel to the $b$ axis, with altern-

[^2]ate layers inverted (Figure 3). The layers are staggered such that the tin(II) lone pair is directed towards a 'hole ' in the network of the next layer (Figure 1).

Bond-distance and -angle data for the two sets of oxygen co-ordination about each tin atom are summarised and compared in Figure 4. Both have pseudo-trigonal-bipyramidal co-ordination with one vacant equatorial position, presumably occupied by the lone pair of electrons of the $\operatorname{tin}$ (II) atom. The $\mathrm{Sn}^{-} \mathrm{O}_{\mathrm{ax}}$. bond distances of each set are identical $[2.36(1) \AA]$, and are comparable with the equivalent bond distances in the $\left[\mathrm{Sn}\left(\mathrm{O}_{4} \mathrm{C}_{2}\right)_{2}\right]^{2-}$ anion ${ }^{8}$ and $\mathrm{Sn}\left(\mathrm{SO}_{4}\right) \cdot 2 \mathrm{SC}\left(\mathrm{NH}_{2}\right)_{2}{ }^{15}$ (Table 4). The shortest such bond distance (except that in $\mathrm{SnO})^{17}$ is observed in $\mathrm{Sn}(\mathrm{PhCO} \cdot \mathrm{CH} \cdot \mathrm{COMe})_{2}[2.290(6)$


Figure 3 Projection on to the $b c$ plane illustrating the layering of the two-dimensional sheets within the crystal
$\AA],{ }^{13}$ whilst the longest occurs in $\mathrm{SnCl}_{2}$-dioxan [2.527(5) $\AA] .^{16}$ The $\mathrm{Sn}-\mathrm{O}_{\text {eq }}$. bond distances are, not unexpectedly, significantly shorter, but are different in the two sets $[2.14(1)$ and $2.20(1) \AA]$. Both, however, are not atypical (Table 4). The $\mathrm{O}_{\text {eq. }}-\mathrm{Sn}-\mathrm{O}_{\text {eq. }}$ and $\mathrm{O}_{\mathrm{ax} .}-\mathrm{Sn}-\mathrm{O}_{\mathrm{ax}}$. bond angles in each set are similar, and exhibit the usual reduction from the values expected for ideal trigonalbipyramidal geometry due to lone-pair-bond-pair repulsion. The $\mathrm{O}_{\mathrm{ax} .}-\mathrm{Sn}-\mathrm{O}_{\mathrm{ax} .}$ angles are not reduced as much as in $\mathrm{Sn}(\mathrm{PhCO} \cdot \mathrm{CH} \cdot \mathrm{COMe})_{2}{ }^{13}$ or in the $\left[\mathrm{Sn}\left(\mathrm{O}_{4} \mathrm{C}_{2}\right)_{2}\right]^{2-}$ anion, ${ }^{8}$ where the ligands chelate the tin, but the restrictions of the two-dimensional network do not allow the bond angle



Figure 4 Comparison of bond-angle and -distance data for the two disordered sets of oxygen co-ordination about the tin
to open to the value observed in the one-dimensional coordinate polymer $\mathrm{SnCl}_{2} \cdot$ dioxan $\left[168.5(2)^{\circ}\right],{ }^{16}$ the highest such bond observed. Viewing down the axis bisecting
the $\mathrm{O}_{\mathrm{ax} .}-\mathrm{Sn}-\mathrm{O}_{\mathrm{ax} .}$. and $\mathrm{O}_{\text {eq. }}-\mathrm{Sn}-\mathrm{O}_{\text {eq. }}$. bond angles of each individual tin atom, the two $\mathrm{SnO}_{4}$ skeletons are mutually rotated by ca. $54^{\circ}$ with respect to each other.

All the formate groups span axial and equatorial coordination sites on adjacent tin atoms, and are of two types. Half the formate groups are symmetrical with a $\mathrm{C}-\mathrm{O}$ bond distance of $1.26(2) \AA$ and an $\mathrm{O}-\mathrm{C}-\mathrm{O}$ bond angle
those where the oxygen bonds to tin. Bridging carboxylate groups occur in $\left[\mathrm{Sn}^{\mathrm{II}} \mathrm{Sn}^{\mathrm{IV}}\left(\mathrm{O}_{2} \mathrm{CC}_{6} \mathrm{H}_{4} \mathrm{NO}_{2}-o\right)_{4} \mathrm{O} \cdot \text { thf] }\right]_{2}{ }^{6}$ with the $\mathrm{C}-\mathrm{O}-\left(\mathrm{Sn}^{\mathrm{II}}\right)$ bonds shorter $[$ mean $1.22(1) \AA]$ than the $\mathrm{C}-\mathrm{O}-\left(\mathrm{Sn}^{\text {IV }}\right)$ distances $[1.26(2)-1.30(1) \AA]$, and the tin-oxygen bond distances are consistent with a $\mathrm{Sn}^{\mathrm{IV}}$ $\left.\mathrm{O}-\mathrm{C}=\mathrm{O} \rightarrow \mathrm{Sn}^{\mathrm{II}}\right]$ electronic arrangement. Thus, in the present compound, two of the formate groups have a

Table 4
Comparison of structural data in $\operatorname{tin}(\mathrm{II})$ compounds possessing a distorted pseudo-trigonal-bipyramidal geometry

|  | Distance/ $\AA$ |  | Angle ( ${ }^{\circ}$ ) |  |
| :---: | :---: | :---: | :---: | :---: |
| Compound | Sn-Oeq. | $\mathrm{Sn}-\mathrm{O}_{\mathbf{a x}}$. |  | $\mathrm{O}_{\mathrm{ax} .}-\mathrm{Sn}-\mathrm{O}_{\mathrm{ax}}$ |
| $\mathrm{Sn}\left(\mathrm{O}_{2} \mathrm{CH}\right)_{2}{ }^{\text {a }}$ (i) | 2.14 (1) | 2.36(1) | 80.7(4) | 158.6(4) |
| (ii) | 2.20 (1) | 2.36(1) | 78.7(4) | 157.7(4) |
| $\mathrm{Sn}(\mathrm{PhCO} \cdot \mathrm{CH} \cdot \mathrm{COMe})_{2}{ }^{\text {b }}$ | $2.135(1)$ | 2.290 (6) | 94.7(1) | 150.4(2) |
| $\left[\mathrm{Sn}\left(\mathrm{O}_{4} \mathrm{C}_{2}\right)_{2}\right]^{2-c}$ | 2.25 (1) | 2.36 (1) | 78.7(3) | 145.5(4) |
| $\mathrm{SnCl}_{2} \cdot{ }^{\text {dioxan }}{ }^{\text {d }}$ |  | 2.527(5) |  | 168.5(2) |
| $\mathrm{Sn}\left(\mathrm{SO}_{4}\right) \cdot 2 \mathrm{SC}\left(\mathrm{NH}_{2}\right)_{2}{ }^{\text {e }}$ | 2.24(2) | 2.41(2) | $89.6(5){ }^{f}$ | $157.3(4){ }^{\text {g }}$ |
| $\mathrm{SnO}{ }^{4}$ | 2.21 (1) | 2.21 | 75(4) | 118(2) |

of $122.6(12)^{\circ}$. The remaining formate groups are unsymmetrical, with one short $[1.23(2) \AA]$ and one long $[1.28(2) \AA]$ bond, although the $\mathrm{O}-\mathrm{C}-\mathrm{O}$ bond angle remains the same $\left[122.3_{5}(13)\right.$ and $\left.123.7(12)^{\circ}\right]$. In the latter case, the oxygen atom involved in the longer bond to carbon is bonded to tin via an equatorial site, and the oxygen atom involved in the longer bond to carbon bonds
greater contribution of the ester $-\mathrm{O}-\mathrm{CH}=\mathrm{O}$ form than the other two, which are more ionic in character. The $\mathrm{Sn}-\mathrm{O}_{\text {eq }}$. -C bond angle varies only slightly [115.7(4)$\left.118.7(4)^{\circ}\right]$; the $\mathrm{Sn}-\mathrm{O}_{\mathrm{ax} .}{ }^{-\mathrm{C}}$ angles are significantly larger [127.5(9)—131.1(9) $\left.{ }^{\circ}\right]$.

The two types of formate groups are manifest in the vibrational spectra, which are listed together with

Table 5
Infrared and Raman bands $\left(\mathrm{cm}^{-1}\right)$ of $\operatorname{tin}(\mathrm{II})$ formate

| Band | Mode |
| :---: | :---: |
| $\nu_{1}$ | $\nu(\mathrm{C}-\mathrm{H})$ |
| $\nu_{4}$ | $\nu_{\text {asym }}(\mathrm{OCO})$ |
| $\nu_{5}$ | $\delta(\mathrm{C}-\mathrm{H})$ (in plane) |
| $\nu_{2}$ | $\nu_{\text {sym }}(\mathrm{OCO})$ |
| $\nu_{6}$ | $\delta(\mathrm{C}-\mathrm{H})$ (out of plane) |
| $\nu_{3}$ | $\begin{aligned} & \delta_{\mathrm{sym}}(\mathrm{OCO}(\mathrm{OCO} \\ & \nu\left(\mathrm{Sn}-\mathrm{O}_{\mathrm{eq}}\right) \\ & \nu\left(\mathrm{Sn}-\mathrm{O}_{\mathrm{eq}}\right) \\ & \nu\left(\mathrm{Sn}-\mathrm{O}_{\mathrm{ax}} .\right) \end{aligned}$ |
|  | $\delta(\mathrm{OSnO})$ |


| $\overbrace{2 v}^{c}$ | $C_{s}^{c}$ |
| :--- | :--- |
| $A_{1}$ | $A$ |
| $B_{1}$ | $B$ |
| $B_{2}$ | $B$ |
| $A_{1}$ | $A$ |
| $B_{2}$ | $B$ |
| $A_{1}$ | $A$ |
|  |  |
|  |  |
|  |  |
|  |  |


|  | I.r. |
| :---: | :---: |
| From ref. 11 | \% $I^{\text {b }}$ |
| 2874 | 2876 |
| 1563 | $1561 \quad 100\}$ |
|  | 1548 100\} |
| $\{1385$ | 1386 26 |
| \{1370 | 1371 39 |
| $1339\}$ | 1341 49 |
| 1311 , | $1330-1320$ (sh) vbr $\}$ |
| 1064 | 1071 1\} |
|  | 1057 1 |
| 781 | 782 31 |
|  | 350-360 (sh) |
|  | 32911 |
|  | 260 9 |
|  | $\{189$ |
|  | $\{111$ |


| Raman |  |
| :---: | :---: |
| กี | $I{ }^{\text {b }}$ |
| 2873 | 66 |
| 1565 | 66 |
| 1388 | 49 ) |
| 1369 | $100\}$ |
| 1349 | 46 |
| 1327 | 63 ) |
| 1069 | $9\}$ |
| 1057 | 9 9 |
| 785 | 17 |
| 345 |  |
| 327 |  |
| 247 |  |

I.r. overtone: $2920\left(1371+1548=2919 \mathrm{~cm}^{-1}\right)$
${ }^{a}$ Symmetry types based on 'local' symmetry for the formate groups. ${ }^{b}$ Relative intensity.
to tin via an axial site. The formate groups in sodium formate ${ }^{18}$ are symmetrical (C-O $1.27 \AA, \mathrm{O}-\mathrm{C}-\mathrm{O} 124^{\circ}$ ), whilst those of ammonium formate ${ }^{\mathbf{1 9}}$ are nearly so ( $\mathrm{C}-\mathrm{O} 1.237$ and $1.246 \AA, \mathrm{O}-\mathrm{C}-\mathrm{O} 126.3^{\circ}$ ), although the network of hydrogen bonds most probably affects the bonding somewhat. In the $\left[\mathrm{Sn}\left(\mathrm{O}_{2} \mathrm{CH}\right)_{3}\right]^{-}$anion ${ }^{7}$ all the formate groups are unidentate, but one is symmetrical $\left[\mathrm{C}-\mathrm{O} 1.24_{5}(2) \AA, \mathrm{O}-\mathrm{C}-\mathrm{O} 120.8(16)^{\circ}\right]$ whilst in the other two groups the $\mathrm{C}-\mathrm{O}$ bond distances differ by $0.1 \AA[\mathrm{C}-\mathrm{O}$ $1.20(2)$ and $1.30(2) \AA, \mathrm{O}-\mathrm{C}-\mathrm{O} 124.3(15)$ and $\left.126.5(15)^{\circ}\right]$. The $\left[\mathrm{Sn}\left(\mathrm{O}_{4} \mathrm{C}_{2}\right)_{2}\right]^{2-}$ anion ${ }^{8}$ has $\mathrm{C}-\mathrm{O}$ bond distances in the range $1.24(2)-1.29(2) ~ \AA$ although the longest bonds are

[^3]assignments in Table 5 and with the data of Donaldson et al. ${ }^{11}$ The majority of the fundamentals of the formate groups have two components, although the assignment of bands to the symmetrical formate groups ( $C_{2 v}$ ' local ' symmetry) or the unsymmetrical formate groups ( $C_{s}$ ' local ' symmetry) is to a large part indeterminate. The bands at $<500 \mathrm{~cm}^{-1}$ are all metal-ligand vibrations. Those at $350-360$ and $329 \mathrm{~cm}^{-1}$ may be assigned to $\mathrm{Sn}^{-}$ $\mathrm{O}_{\text {eq. }}$ stretching vibrations, the higher-energy band to the shorter bond, and vice versa. The $\mathrm{Sn}-\mathrm{O}_{\mathrm{ax}}$. stretching vibration is assigned to the band at $260 \mathrm{~cm}^{-1}$. Bands at 189 and $111 \mathrm{~cm}^{-1}$ are assigned as $\mathrm{O}^{-} \mathrm{Sn}^{-} \mathrm{O}$ angle-deformation modes.

The tin-119m Mössbauer spectrum of $\operatorname{tin}$ (II) formate
has been measured previously at 77 K by Donaldson and Jelen, ${ }^{3}$ and the sign of the quadrupole coupling constant, $e V_{z z} Q$, has been shown to be positive. ${ }^{10}$ As expected,

Table 6
Tin-119m Mössbauer data ( $\mathrm{mm} \mathrm{s}^{-1}$ ) for $\operatorname{tin}(\mathrm{II})$ formate ${ }^{a}$

| $T / \mathrm{K}$ | I.s. | Q.s. | $\Gamma_{1}$ | $\Gamma_{2}$ | $A_{1} / A_{2}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| 77 | 3.33 | 11.70 | 1.02 | 1.11 | 0.78 |
| Ambient | 3.27 | 1.61 | 0.77 | 0.78 | 0.74 |
| $b$ | 3.15 | 1.56 |  |  |  |

${ }^{a}$ Relative to $\mathrm{BaSnO}_{3}\left(0 \mathrm{~mm} \mathrm{~s}^{-1}\right)$. ${ }^{b}$ Data at 77 K from ref. 3. because of its polymeric lattice, tin(II) formate exhibits a Mössbauer spectrum at both 77 K and ambient tem-
perature. The data are listed in Table 6. Both the isomer shift and quadrupole splitting decrease with increasing temperature, and are somewhat higher than those reported previously. A distinct GoldaskiiKaryagin effect is observed at both temperatures. A more detailed investigation of the temperature dependence of the Mössbauer parameters will be reported subsequently.

We thank the S.R.C. for support, and Professor T. J. King for collecting the $X$-ray intensity data and helpful discussions.
[7/1973 Received, 9th November, 1977]


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